

# Symbolic Verification of Quantum Protocols via Quantum Distributions – Early Ideas

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## Abstract

The development of quantum communication protocols sparked the interest in quantum extensions of process calculi. The prominent approaches, however, are still based on probabilistic models, which are inadequate for the quantum case, as they do not capture the observational limitations of quantum theory. In Physics, probabilistic quantum systems are not represented with probabilities, but with density operators. By lifting this concept to the Computer Science setting, we introduce Operator Distributions: they employ density operators as weights, just like probability distributions use real numbers. We give semantics to quantum protocols in a natural way, as an LTS of operator distributions, guaranteeing that only physically accessible information about the systems are represented. Moreover, these “quantum” distributions pave the way for symbolic verification of protocols whose behaviour depends on an unknown initial state: instead of considering infinitely-many LTSs, one for each possible initialization, we can build a single symbolic LTS and prove properties on it.

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Recent years have seen a flourishing development of *quantum computation* and *quantum communication* technologies [22]. We focus on the latter, which provides solutions for key distribution [23], electronic voting [2], *Quantum Internet* [3, 26], and many other applications [1]. Therefore, the need emerged for modelling and verification techniques applicable to quantum protocols. We will show why the current proposals are not fully satisfactory, and then present two novel semantic models based on *Quantum Distributions*, which abides the rules of quantum theory and supports algorithmic equivalence checking. This is useful to prove correctness of protocol implementations, or to investigate non-interference and resistance to cryptographic attackers.

Numerous works in the field [21, 17, 15, 7, 6] rely on *quantum process calculi*, where we describe a set of interacting agents, each of them manipulating and sending some qubits. The typical semantics for these quantum calculi is based on *configurations*  $\langle |\psi\rangle, P \rangle$ , consisting of a syntactic *process*  $P$  (the code that the agent is executing) and a quantum state  $|\psi\rangle$  (the data which the process manipulates). An example would be the configuration  $C = \langle |0\rangle, H(q).M_{01}(q \triangleright x).c!q \rangle$ , where the data is a single qubit  $q$  in state  $|0\rangle$ , and the process  $i$ ) updates  $q$  with an  $H$  gate, bringing its state to  $|+\rangle$ ; ii) measures such qubit and stores the outcome in the variable  $x$ , causing  $q$  do *decay* in either  $|0\rangle$  or  $|1\rangle$ ; iii) sends  $q$  over the channel  $c$ . In literature, it is common to build a Labelled Transition System (LTS) whose states are *distributions* of such configurations. The LTS of the distribution  $1 \cdot C$  would be

$$1 \cdot C \xrightarrow{\tau} 1 \cdot \langle |+\rangle, \underline{M_{01}(q \triangleright x).c!q} \rangle \xrightarrow{\tau} \frac{1}{2} \cdot \langle |0\rangle, \underline{c!q} \rangle \oplus \frac{1}{2} \cdot \langle |1\rangle, \underline{c!q} \rangle \xrightarrow{c!q} \frac{1}{2} \cdot \langle |0\rangle, \underline{\mathbf{0}} \rangle \oplus \frac{1}{2} \cdot \langle |1\rangle, \underline{\mathbf{0}} \rangle \quad (1)$$

47 where  $1 \cdot C$  is the point distribution in  $C$ , and  $\frac{1}{2} \cdot C \oplus \frac{1}{2} \cdot C'$  is the uniform distribution of  $C, C'$ .  
 48 Notice how configurations in quantum process calculi are akin to “quantum closures” in the  
 49 operational semantics of quantum languages [25]. The approach of modelling probabilistic  
 50 systems with a transition systems of distributions has already appeared for both concurrent  
 51 systems [19, 11] and quantum lambda calculi [8, 12, 13].

52 Our main focus in characterizing when two protocols are *behaviourally equivalent*, meaning  
 53 that they exhibit the same interactive observable behaviour, in term of communication actions  
 54 and qubit values. Different notions of behavioural equivalence have been considered [9, 16, 10,  
 55 5], all based on the distribution-of-configurations approach described above. However argue  
 56 that it is the semantic model itself to be inadequate, for two reasons: it does not respect  
 57 the indistinguishability results of quantum theory, and it hampers algorithmic equivalence  
 58 checking.

59 For the first problem, consider these two sources of random qubits:

$$60 \quad \Delta = \frac{1}{2} \cdot \langle |0\rangle, \underline{c!q} \rangle \oplus \frac{1}{2} \cdot \langle |1\rangle, \underline{c!q} \rangle \quad \Theta = \frac{1}{2} \cdot \langle |+\rangle, \underline{c!q} \rangle \oplus \frac{1}{2} \cdot \langle |-\rangle, \underline{c!q} \rangle. \quad (2)$$

61 The  $\Delta$  distribution sends a random value between  $|0\rangle$  or  $|1\rangle$ , while  $\Theta$  sends either  $|+\rangle$  or  $|-\rangle$ .  
 62 According to quantum mechanics, no realizable procedure can tell apart one source from the  
 63 other [22]. Nonetheless, several existing proposals fail to make  $\Delta$  and  $\Theta$  equivalent [20, 15, 11],  
 64 while others have to fine tune their notion of equivalence to solve this issue [10, 5, 6]. The  
 65 second problem, instead, involves protocols featuring quantum inputs (e.g. Teleportation):  
 66 to check their correctness we should build an infinite number of LTSs, one for each input.

67 We address both concerns by introducing two novel semantic models, inspired by quantum  
 68 theory. In Physics, when dealing with random qubits, we do not use probability distributions  
 69 of quantum values, but we employ the *Density Operator Formalism*. Qubits states are repres-  
 70 ented as matrices, called *Density Operators*, and transformations on qubits are represented as  
 71 maps between matrices, called *Superoperators*. The density operators on  $d$ -dimensional Hilbert  
 72 Space  $\mathcal{H}$  are denoted as  $DO_{\mathcal{H}} = \{\rho \in \mathbb{C}^{d \times d} \mid \rho \text{ is positive and } \text{tr}(\rho) \leq 1\}$ , while superoperat-  
 73 ors are defined as  $Sop_{\mathcal{H}} = \{\mathcal{E} : DO_{\mathcal{H}} \rightarrow DO_{\mathcal{H}} \mid \mathcal{E} \text{ is completely positive and } \text{tr}(\mathcal{E}(\rho)) \leq \text{tr}(\rho)\}$ .

74 There is a non-injective map from probability distributions of quantum values to density  
 75 operators. Distributions that are impossible to tell apart are mapped in the same matrix.  
 76 For example, the two distribution  $\frac{1}{2} \cdot |0\rangle \oplus \frac{1}{2} \cdot |1\rangle$  and  $\frac{1}{2} \cdot |+\rangle \oplus \frac{1}{2} \cdot |-\rangle$  are both represented as

$$77 \quad \frac{1}{2} |0\rangle\langle 0| + \frac{1}{2} |1\rangle\langle 1| = \mathbb{I}/2 = \frac{1}{2} |+\rangle\langle +| + \frac{1}{2} |-\rangle\langle -| \quad (3)$$

78 Density operators enjoy a *Partial Commutative Monoid* (PCM) structure: they can be  
 79 summed with matrix addition, but only if the result is still in  $DO_{\mathcal{H}}$ . This generalizes the  
 80 case of probabilities, i.e. the PCM of real numbers with sum in the interval  $[0, 1]$ . Following  
 81 this similarity, we can define a new kind of *quantum distributions*.

82 ► **Definition 1.** *Given a Hilbert Space  $\mathcal{H}$ , the set of Operator (Sub)-Distributions over  $X$  is*

$$83 \quad \mathcal{Q}(X) = \{ \mathfrak{D} \in (DO_{\mathcal{H}})^X \mid \text{supp}(\mathfrak{D}) \text{ is finite and } \sum_{x \in X} \mathfrak{D}(x) \in DO_{\mathcal{H}} \}$$

84 *where  $\text{supp}(\mathfrak{D})$  is the support of  $\mathfrak{D}$ , i.e. the set  $\{x \in X \mid \mathfrak{D}(x) \neq \mathbf{0}\}$ .*

85 We write  $\rho_0 \cdot x_0 \oplus \rho_1 \cdot x_1$  for the distribution where  $x_i$  has weight  $\rho_i$  for  $i \in \{0, 1\}$ , and all  
 86 other  $x_i \in X$  have weight  $\mathbf{0}$ , the all-zero matrix. Categorically, operator distributions form a  
 87 functor on **Set**, and indeed we can build similar “distribution functors” for any *PCM*.

88 We can give semantics to quantum protocols using *operator distribution of processes*. Each  
 89 process is weighted with an operator, encoding both the quantum data and the probability  
 90 of having reached that process. The LTS in Equation 1 would become

$$91 \quad |0\rangle\langle 0| \cdot \underline{H(q).M_{01}(q \triangleright x).c!q} \xrightarrow{\tau} |+\rangle\langle +| \cdot \underline{M_{01}(q \triangleright x).c!q} \xrightarrow{\tau} \mathbb{I}/2 \cdot \underline{c!q} \xrightarrow{c!q} \mathbb{I}/2 \cdot \underline{\mathbf{0}}. \quad (4)$$

92 Indeed, after measuring  $q$ , the system evolves in the distribution  $\frac{1}{2} |0\rangle\langle 0| \cdot \underline{c!q} \oplus \frac{1}{2} |1\rangle\langle 1| \cdot \underline{c!q}$ . But  
 93  $\frac{1}{2} |0\rangle\langle 0| \cdot \underline{c!q} \oplus \frac{1}{2} |1\rangle\langle 1| \cdot \underline{c!q}$  is just another way to write  $\mathbb{I}/2 \cdot \underline{c!q}$ , thanks to Equation 3. This solves  
 94 the problem in Equation 2: equivalent sources are modelled as the same operator distribution.  
 95 Notice how this would not happen for the process  $\underline{H(q).M_{01}(q \triangleright x).c!x}$ , which measures the  
 96 qubit and sends the classical outcome. It evolves in  $\frac{1}{2} |0\rangle\langle 0| \cdot \underline{c!0} \oplus \frac{1}{2} |1\rangle\langle 1| \cdot \underline{c!1}$ , where the  
 97 processes are *classically* distinguished, and thus the quantum states are not combined.

98 However, the process  $\underline{P} = \underline{H(q).M_{01}(q).c!q}$  in Equation 4 has still a different transition  
 99 system for all its inputs, like  $|0\rangle\langle 0| \cdot \underline{P}$ ,  $|+\rangle\langle +| \cdot \underline{P}$  or  $\mathbb{I}/2 \cdot \underline{P}$ . We thus introduce a second kind  
 100 of quantum distributions, based on superoperators, inspired by the denotational semantics of  
 101 quantum programming languages [24, 18]. Similarly to  $DO_{\mathcal{H}}$  and to  $[0, 1]$ , also superoperators  
 102 on  $\mathcal{H}$  enjoy a PCM structure, such that we can define *Superoperator Distributions* as follows.  
 103

104 **► Definition 2.** *Given a Hilbert Space  $\mathcal{H}$ , we define Superoperator Distributions over  $X$  as*

$$105 \quad \mathcal{S}(X) = \{ \mathfrak{D} \in (\text{Sop}_{\mathcal{H}})^X \mid \text{supp}(\mathfrak{D}) \text{ is finite and } \sum_{x \in X} \mathfrak{D}(x) \in \text{Sop}_{\mathcal{H}} \}$$

106 where  $\text{supp}(\mathfrak{D}) = \{ x \in X \mid \mathfrak{D}(x) \neq \mathcal{Z} \}$ ,  $\mathcal{Z}$  is the constant  $\mathbf{0}$  map, and superoperators are  
 107 summed in a point-wise manner.

108 The semantics of a protocol can then be given as a single LTS made of superoperator  
 109 distributions, similar to [14, 4]. The LTS in Equation 1 becomes

$$110 \quad \underline{\text{Id} \cdot H(q).M_{01}(q \triangleright x).0} \xrightarrow{\tau} \underline{\text{Had} \cdot M_{01}(q \triangleright x).c!q} \xrightarrow{\tau} \underline{\text{Set}_{\mathbb{I}/2} \cdot c!q} \xrightarrow{c!q} \underline{\text{Set}_{\mathbb{I}/2} \cdot \mathbf{0}}. \quad (5)$$

111 where  $\text{Had}$  is the Hadamard gate, and  $\text{Set}_{\mathbb{I}/2}$  is the map which prepares the state  $\mathbb{I}/2$ . In  
 112 these distributions, each process has a superoperator as weight, representing the sequence of  
 113 transformation that have been applied on the quantum input to arrive at that process.

114 Noticeably, this LTS is independent on the state of the input qubit. We can check our  
 115 behavioural equivalence of choice, e.g. bisimilarity, between these *symbolic*, superoperator-  
 116 based transition systems, instead of comparing an infinite number of *ground*, operator-based  
 117 ones. Once we know the input quantum state, each symbolic LTS can be instantiated to a  
 118 ground LTS, just by updating the weights in the distributions. In the case of Equation 5, if  
 119 the input state is  $|0\rangle\langle 0|$  we get the LTS in Equation 4.

120 Categorically, each input state  $\rho$  defines a natural transformation between superoperator  
 121 distributions and operator distributions, and these transformations commute with the  
 122 semantics of processes. From this, it is possible to prove full abstraction: two symbolic LTSs  
 123 are bisimilar if and only if all their ground instantiations are.

124 More in detail, we can lift some useful properties from our weights ( $DO_{\mathcal{H}}$  and  $\text{Sop}_{\mathcal{H}}$ ) to  
 125 our distributions ( $\mathcal{Q}(X)$  and  $\mathcal{S}(X)$ ). We can show that  $\text{Sop}_{\mathcal{H}}$  forms an appropriate kind  
 126 of *partial semiring*, with partial addition and total multiplication, while both  $DO_{\mathcal{H}}$  and  
 127  $\text{Sop}_{\mathcal{H}}$  form a *partial Sop-semimodule*, as they can be summed and can be “updated” by a  
 128 superoperator. Then, given a quantum input  $\rho$ , instantiating a superoperator  $\mathcal{E} \in \text{Sop}_{\mathcal{H}}$   
 129 to get a quantum output  $\mathcal{E}(\rho) \in DO_{\mathcal{H}}$  is a *homomorphism of Sop-semimodules*: updating  
 130 our symbolic weight  $\mathcal{E}$  with another superoperator  $\mathcal{F}$  commutes with instantiating it to a

131 ground weight  $\mathcal{E}(\rho)$ . Lifting these properties from weights to distributions and LTSs, get a  
 132 fully abstract translation between our symbolic  $Sop_{\mathcal{H}}$ -weighted semantics and our ground  
 133  $DO_{\mathcal{H}}$ -weighted one.

134 To sum up, we propose two quantum generalization of probability distributions, suitable  
 135 to model and verify quantum communication protocols. Our ground semantics is defined in  
 136 terms of  $DO_{\mathcal{H}}$ -weighted distributions of processes, where density operators represent both  
 137 quantum data and probabilities, and they naturally abide to the observational limitations  
 138 of quantum theory. Then, to check equivalence between quantum protocols in a symbolic  
 139 fashion, we can define  $Sop_{\mathcal{H}}$ -weighted distributions, employing superoperators to represent  
 140 the sequence of transformations applied on an unknown quantum state. Thanks to full  
 141 abstraction, we can check equivalence between symbolic systems, and know that they will be  
 142 equivalent under any possible quantum input.

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